



### New directions in phase-field modeling of microstructure evolution in polycrystalline and multi-component alloys

Nele Moelans

Liesbeth Vanherpe, Jeroen Heulens, Bert Rodiers

K.U. Leuven, Belgium



### 'Quantitative' phase-field models

- Properties bulk and interfaces are reproduced accurately in the simulations
  - Effect model description and parameters
  - Numerical issues
- Insights in the evolution of complex morphologies and grain assemblies
  - Effect of individual bulk and interface properties
- Predictive ?
  - Depends on availability and accuracy of input data
    - Requires composition and orientation dependence









### General framework and goal







### Some aspects of model formulation

- 2-phase systems (single phase-field)
- Multi-grain/phase systems (multiple phase-field)



### 2-phase systems

• Field variables:  $\phi(\vec{r},t) = c_k(\vec{r},t)$ 







$$g(\phi) = 16\phi^{2}(1-\phi)^{2}$$

$$1 = 16\phi^{2}(1-\phi)^{2}$$

$$0 = 16\phi^{2}(1-\phi)^{2}$$

$$0 = 16\phi^{2}(1-\phi)^{2}$$

Double well function

Free energy  

$$F = \int_{V} \left[ f_{chem}(c,\phi) + \underbrace{Wg(\phi) + \frac{\varepsilon^{2}}{2} |\nabla\phi|^{2}}_{Interfacial \ energy} \right] d\vec{r}$$

• Bulk energy  $f_{chem} = h(\phi) f^{\beta}(c,T) + \left[1 - h(\phi)\right] f^{\alpha}(c,T)$ 





# **LEUVEN Decoupling bulk and interfacial energy**

- Interface treated as mixture of 2 phases
  - c-field for each phase  $c \rightarrow c^{\alpha}, c^{\beta}$
  - Equal interdiffusion potential + conservation

$$\begin{cases} \frac{\partial f^{\beta}(c^{\beta})}{\partial c^{\beta}} = \frac{\partial f^{\alpha}(c^{\alpha})}{\partial c^{\alpha}} = \tilde{\mu}\\ c = h(\phi)c^{\beta} + \left\lceil 1 - h(\phi) \right\rceil c^{\alpha} \end{cases}$$





Kim et al., PRE, 6 (1999) p 7186; Tiaden et al., Physica D, 115 (1998) p73

• Bulk energy

$$\Rightarrow f_{chem} = h(\phi) f^{\beta}(c^{\beta}) + \left[1 - h(\phi)\right] f^{\alpha}(c^{\alpha})$$



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# Decoupling bulk and interfacial kinetics

- Kinetic equations (Linear non-equilibrium thermodynamics)
  - Allen-Cahn  $\frac{\partial \varphi}{\partial r} = -M_{\varphi} \frac{\partial F}{\partial r}$

Diffusion

$$\frac{\partial t}{\partial t} = \nabla \cdot \sum_{l=1}^{C-1} \left[ h(\phi) M_{kl}^{\beta} + [1 - h(\phi)] M_{kl}^{\alpha} \right] \nabla \tilde{\mu}_{l}$$

Jump in chemical potential accross interface



$$\frac{\partial c_{k}}{\partial t} = \nabla \cdot [1 - h(\phi)] \sum_{l=1}^{C-1} M_{kl}^{L} \nabla \tilde{\mu}_{l} + \nabla \cdot \alpha_{i} \frac{\partial \phi}{|\partial t| |\nabla \phi|} \frac{\nabla \phi}{|\nabla \phi|}$$
*Non-variational anti-trapping current*

Dilute, D<sub>S</sub>=0: A.Karma, PRL, 87, 115701 (2001); B. Echebarria et al., PRE, 70, 061604 (2004)
 Multi-comp, D<sub>S</sub>=0: S.G. Kim, Acta Mater. 55, p4391 (2007)



### Multi-grain and multi-phase structures

 Single phase-field models -> Multiple phase-field models

 $\eta \rightarrow \left\{ \eta_1, \eta_2, \eta_3, ..., \eta_p \right\}$ 

 $(\eta_1, \eta_2, ..., \eta_i, ..., \eta_p) = (0, 0, ..., 1, ..., 0)$ 

Model extension

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 $F(\eta_1, \eta_2, \eta_3, ..., |\nabla \eta_1|^2, |\nabla \eta_2|^2, ...)$ 

- Different types of interfaces
- Triple and higher order junctions
- Numerically
  - Same accuracy for all interfaces and phases
  - All interfaces within range of validity of the thin interface asymptotics

$$\rightarrow \ell_{num} = cte$$









### Multi-grain and multi-phase models: major difficulties

- Third-phase contributions
  - $σ_{12} = σ_{13} = 7/10 σ_{12}$



- → Careful choice of multi-well function and gradient contribution
- Interpolation function
  - Zero-slope at equilibrium values of the phase fields
  - Thermodynamic consistency

$$f_{chem} = \sum_{i=1}^{p} h_i(\eta_1, \eta_2, ...) f^i(c, T) \implies \sum_{i=1}^{p} h_i(\eta_1, \eta_2, ...) = 1$$









### Anisotropic grain growth model

- Phase fields  $\eta_1, \eta_2, ..., \eta_i(\vec{r}, t), ..., \eta_p$ 
  - With grain i

 $(\eta_1, \eta_2, ..., \eta_i, ..., \eta_p) = (0, 0, ..., 1, ..., 0)$ 

• Free energy

$$F_{interf} = \int_{V} m \left( \sum_{i=1}^{p} \left( \frac{\eta_i^4}{4} - \frac{\eta_i^2}{2} \right) + \sum_{i=1}^{p} \sum_{j < i}^{p} \gamma_{i,j} \eta_i^2 \eta_j^2 + \frac{1}{4} \right) + \frac{\kappa(\eta)}{2} \sum_{i=1}^{p} (\nabla \eta_i)^2 dV$$



$$\boldsymbol{\kappa}(\boldsymbol{\eta}) = \sum_{i=1}^{p} \sum_{j < i}^{p} \boldsymbol{\kappa}_{i,j} \boldsymbol{\eta}_{i}^{2} \boldsymbol{\eta}_{j}^{2} / \sum_{i=1}^{p} \sum_{j < i}^{p} \boldsymbol{\eta}_{i}^{2} \boldsymbol{\eta}_{j}^{2}$$

• For each grain boundary 
$$\eta_i^2 \eta_j^2 \neq 0 \Rightarrow \kappa(\eta) = \kappa_{i,j}$$

Inclination dependence

$$\gamma_{i,j}\left(\boldsymbol{\psi}_{i,j}\right), \boldsymbol{\kappa}_{i,j}\left(\boldsymbol{\psi}_{i,j}\right), \boldsymbol{L}_{i,j}\left(\boldsymbol{\psi}_{i,j}\right), \qquad \boldsymbol{\psi}_{i,j} = \frac{\nabla \boldsymbol{\eta}_{i} - \nabla \boldsymbol{\eta}_{j}}{|\nabla \boldsymbol{\eta}_{i} - \nabla \boldsymbol{\eta}_{j}|}$$



L.-Q. Chen and W. Yang, PRB, 50 (1994) p15752 A. Kazaryan et al., PRB, 61 (2000) p14275

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### Non-variational approach – equal interface width

Ginzburg-Landau type equations

$$\frac{\partial \eta_i(\vec{r},t)}{\partial t} = -L(\eta) \left[ m \left( \eta_i^3 - \eta_i + 2\eta_i \sum_{j \neq i} \gamma_{i,j} \eta_j^2 \right) - \kappa(\eta) \nabla^2 \eta_i \right]$$

- Non-variational with respect to η-dependence of κ
- Similar to Monte Carlo Potts approach
- Definition 'grain boundary width'

$$l_{num} = \frac{1}{\left|\frac{d\eta_i}{dx}\right|_{max}} = \frac{1}{\left|\frac{d\eta_j}{dx}\right|_{max}}$$

→High controllability of numerical accuracy ( $I_{num}/R < 5$ )







### **Grain boundary properties**

Grain boundary energy

$$\gamma_{gb,\theta_{i,j}} = g(\gamma_{i,j}) \sqrt{m\kappa_{i,j}}$$

Grain boundary mobility

 $\mu_{gb,\theta_{i,j}} = L_{i,j} \sqrt{\frac{\kappa_{i,j}}{m(g(\gamma_{i,j}))^2}}$ 

Grain boundary width

$$l = \frac{4}{3} \sqrt{\frac{\kappa_{i,j}}{m(g(\gamma_{i,j}))^2}}$$



 $g(\gamma_{i,i})$  calculated numerically

• Iterative algorithm

$$\ell_{gb}, [\gamma_{gb,\theta}], [\mu_{gb,\theta}] \to m, [\kappa_{i,j}], [\gamma_{i,j}], [L_{i,j}]$$

N. Moelans, B. Blanpain, P. Wollants, PRL, 101, 0025502 (2008); PRB, 78, 024113 (2008)





### **Numerical validation**



- Diffuse interface effects for  $\ell_{num}$  / R > 5
- Angles outside [100°-140°] require larger  $\ell_{num}$  /  $\Delta x$  for same accuracy

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### Variational approach – interface width varies with interface energy

1 r P

- Anisotropy only through  $\gamma_{i,p}$   $\kappa$ =cte 0
  - Energy •
  - **Kinetic equations**  $\odot$

$$F = \int_{V} \left[ m \left( \sum_{i=1}^{p} \left( \frac{\eta_{i}^{4}}{4} - \frac{\eta_{i}^{2}}{2} \right) + \sum_{i=1}^{p} \sum_{j < i}^{p} \gamma_{i,j} \eta_{i}^{2} \eta_{j}^{2} + \frac{1}{4} \right) + \frac{\kappa}{2} \sum_{i=1}^{p} (\vec{\nabla} \eta_{i})^{2} \right] dV$$
$$\frac{\partial \eta_{i}(\vec{r}, t)}{\partial t} = -L(\eta) \left[ m \left( \eta_{i}^{3} - \eta_{i} + 2\eta_{i} \sum_{j \neq i} \gamma_{i,j} \eta_{j}^{2} \right) - \kappa \nabla^{2} \eta_{i} \right]$$







0

### Rotation invariance of the model

- Mathematically, the model equations are invariant to rotation, but ...
- the order parameters represent <u>orientations</u> in a fixed reference frame.



• For the model to be rotational invariant in practice, lower limit of amount of order parameters p:  $\sqrt{2}\pi L$ 

$$p > \frac{\sqrt{2}\pi L}{n\Delta h}$$

J. Heulens and N. Moelans, Scripta Mat. (2010)





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### Extension to multi-component multiphase alloys

• Phase field variables:

Composition

• Grains

$$\eta_{\alpha 1}, \eta_{\alpha 2}, ..., \eta_{\alpha i}(\vec{r}, t), ..., \\\eta_{\beta 1}, \eta_{\beta 2}, ..., \eta_{p} \\\vec{c}_{A}, \vec{c}_{B}(\vec{r}, t), ..., \vec{c}_{C-1}$$

• Bulk energy: 
$$f_{bulk}(c_k,\eta_{\rho i}) = \sum_{\rho} \phi_{\rho} f^{\rho}(c_k^{\rho})$$

• with







Nele Moelans Third annual workshop HERO-M, Saltsjöbaden, Sweden, May 17-18, 2010  $\sum \eta_{
ho i}^2$ 

 $\sum \eta_{\pi i}^2$ 





### Extension to multi-component multiphase alloys

• Bulk and interface diffusion:

Vith 
$$M_k^{\rho} = \frac{D_k^{\rho}}{\frac{\partial^2 f_m^{\rho}}{\partial x_k^2}}$$
 and

$$\frac{\partial x_k}{\partial t} = \nabla \cdot \left[ \left( \sum_{\rho} \phi_{\rho} M_k^{\rho} + \sum_{\rho, i \neq \sigma, j} \eta_{\rho, i}^2 \eta_{\sigma, j}^2 \right) \nabla \mu_k \right]$$

$$M_{interf} = 3 \left( \frac{D_{interf}}{\partial^2 f^m / \partial x_k^2} \right) \left( \frac{\delta_{gb}}{\delta_{num}} \right)$$

• Interface movement:

$$\frac{\partial \eta_{i\rho}}{\partial t} = -L \frac{\delta F(\eta_{i\rho}, x_k)}{\delta \eta_{i\rho}}$$

• Between phase  $\alpha$  and  $\beta$ 

$$\frac{\partial \eta_{\alpha i}}{\partial t} = -L \left( \underbrace{2\eta_{\alpha i}\eta_{\beta j}^{2}}_{(\eta,\nabla\eta)} + \underbrace{2\eta_{\alpha i}\eta_{\beta j}^{2}}_{(\eta_{\alpha}^{2} + \eta_{\beta}^{2})^{2}} \underbrace{f^{\alpha}(c^{\alpha}) - f^{\beta}(c^{\beta}) - (c^{\alpha} - c^{\beta})\mu}_{Bulk \ energy \ driven} \right)$$
Curvature driven
$$Bulk \ energy \ driven$$





### Numerical validation for multicomponent multi-phase model







### **Bounding Box Implementation**

### • Basic elements

- A grain is set of connected grid points r where |eta\_i(r)| > epsilon
- For each grain, the corresponding bounding box is the smallest cuboid containing the grain



### Algorithm

Solve the equations only locally, inside bounding boxes Only values inside boxes are kept in memory Boxes grow or shrink with grain

Object Oriented C++ implementation

In collaboration with L. Vanherpe and S. Vandewalle, K.U. Leuven (*Vanherpe et al., PRE, 76, n°056702 (2007)*)







### **Application examples**

 Grain growth in anisotropic systems with a fiber texture

*In collaboration with* F. Spaepen, Harvard University



# Grain growth in columnar films with fiber texture

- Grain boundary energy:
  - Fourfold symmetry
  - Extra cusp at  $\theta$  = 37.5°
  - Read-shockley



Discrete orientations

 $\eta_1, \eta_2, ..., \eta_i(r, t), ..., \eta_{60} \Rightarrow \Delta \theta = 1.5^\circ$ 

- Constant mobility
- Initially random grain orientation and grain boundary type distributions



### 2D simulation

*White:*  $\theta = 1.5$  *Gray:*  $\theta = 3$  *Red:*  $\theta = 37,5$ *Black:*  $\theta > 3, \theta \neq 37.5$ 





### Simulations: 1 high-angle energy cusp



- High-angle grain boundaries form independent network
- Low-angle grain boundaries follow movement of high-angle grain boundaries → elongate
- No stable quadruple junctions

White:  $\theta = 1.5$ Gray:  $\theta = 3$ Black:  $\theta > 3$ ,  $\theta \neq 37.5$ Red:  $\theta = 37,5$ 





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### Misorientation distribution function (MDF)

#### Area weigthed MDF --t=0.25(s) θ=1.5° 0.15 -•--t=90 ·θ=37.5° 0.15 ----t=250 θ=3° Length density 2000 Length density ▼ t=418 θ**=36°** t=490 θ**=39°** 0.1 θ**=4.5°** 0.05

40

Read-Shockley + cusp at  $\theta$  = 37.5 °

200

Time (s)

100

0

300

Reaches a stead-state

20

Misorientation  $\theta$  (°)

30

• Low energy boundaries lengthen + their number increases

10

0L 0



500

400







• Grain growth exponent

 PFM: steady-state growth with
 n ≈ 1

- Previous findings: n = 0.6...1
- Mean field analysis:

$$n_{eff} = \frac{kt}{A_h(0) + kt} \to 1, t \to \infty$$
$$k_{eff} = \frac{dA_{eff}}{dt} = \frac{k}{1 + N ? N}$$

Read-Shockley ( $\theta_m$ =15°) + cusp at  $\theta$ =37.5°

N. Moelans, F. Spaepen, P. Wollants, Phil. Mag., 90 p 501-523 (2010)







### **Application examples**

 Coarsening of Al<sub>6</sub>Mn precipitates located on a recrystallization front in Al-Mn alloys

*In collaboration with* A. Miroux, E. Anselmino, S. van der Zwaag, T. U. Delft



### Jerky motion during recrystallization in Al-Mn alloy

- In-situ EBSD observation of recrystallization in AA3103 at 400 °C
  - CamScan X500 Crystal Probe FEGSEM
- Jerky grain boundary motion
  - Stopping time: 15-25 s
  - Pinning by second-phase precipitates
    - $AI_6(Fe,Mn)$ ,  $\alpha$ - $AI_{12}(Fe,Mn)_3Si$
- Added to phase field model
  - Grain boundary diffusion
  - Driving force for recrystallization



**20**µm





### Phase field model

• Multiple order parameter  $\eta_{m,1}(\vec{r},t), \eta_{m,2}(\vec{r},t), ..., \eta_{p,i}(\vec{r},t), ...$ representation:

 $(\eta_{m,1},\eta_{m,2},...,\eta_{p,i},...) = (1,0,...,0,...), (0,1,...,0,...), ...(0,0,...,1,...), ...$ 

- Mn composition field:  $x_{Mn}(r,t)$
- Homogeneous driving pressure for recrystallization: m<sub>d</sub>
- Bulk diffusion + Surface diffusion







### Material properties at 723K

Grain boundary energy high angle	$\gamma_h = 0.324 \text{ J/m}^2$
Interfacial energy Al <sub>6</sub> Mn precipitates	$\gamma_{\rm pr}=0.3~J/m^2$
Mobility high angle grain boundary	$M_{h} = 2.94 \cdot 10^{-11} \text{ m}^2 \text{s/kg}$
At solute content 0.3w% Mn	(Miroux et al.,Mater. Sci. Forum,467-470,393(2004))
Equilibrium composition of matrix	c <sub>Mn,eq</sub> = 0.0524 w% (0.02456 at%)
	(PhD thesis Lok 2005)
Actual composition of matrix	c <sub>Mn</sub> = 0.3 w% (0.1474 at%)
(supersaturated)	(PhD thesis Lok 2005)
Mn diffusion in fcc Al	$D_{0,bulk} = 10^{-2} \text{ m}^2/\text{s}, \text{ Q}_{bulk} = 211 \text{ kJ/mol}$
	→D <sub>bulk</sub> = 5.5973·10 <sup>-18</sup> m²/s
Pipe diffusion high angle boundaries,	$D_{0,p} = D_{0,bulk}, \ Q_p = 0.65 Q_{bulk}$
precipitate/matrix interface	$\rightarrow D_{p} = 1.2195 \cdot 10^{-12} m^{2}/s$
Bulk energy density: $f^{\rho} = A^{\rho}(x-x^{\rho}_{0})^{2}$	$A^m = 6 \cdot 10^{11}; x^m_0 = 0.000258$
	$A^{p} = 6 \cdot 10^{12}; x^{p}_{0} = 0.1429$





# Precipitate coarsening and unpinning

- $P_D < P_{ZS} (P_D \approx P_{ZS})$ 
  - Pinning: P<sub>zs</sub>=3.6 MPa
  - Rex: P<sub>D</sub>=3.1 MPa

 Unpinning mainly through surface diffusion around precipitates









• **P**<sub>D</sub> <<< **P**<sub>ZS</sub>

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- Pinning: P<sub>zs</sub> = 3.6 MPa
- Rex: P<sub>D</sub> = 1.1 MPa

 Unpinning through grain boundary diffusion











### **Application examples**

 Coarsening and diffusion controlled growth in leadfree solder joints

> Within the framework of COST MP-0602 (Advanced Solder Materials for High Temperature Application), Chairs A. Kroupa and A. Watson

# **LEUVEN** Coarsening in Sn(-Ag)-Cu solder joints

- IMC formation and growth precipitate growth Kirkendal voids stresses – grain boundary diffusion
  - CALPHAD description
  - Diffusion coefficients, growth coefficient for IMC-layers



SEM-image of Sn – 3.8Ag–0.7 Cu alloy after annealing for 200h at 150 ℃ (Peng 2007)





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## Phase field model

- **Multiple order parameter model:** 0
  - Grains and phases  $\bullet$

 $\eta_{(Cu),1}, \eta_{(Cu),2}, ..., \eta_{(Cu),i}(r,t), ...,$  $\eta_{Cu_3Sn,1}, \eta_{Cu_3Sn,2}, \dots$  $\eta_{Cu_6Sn_5,1},\eta_{Cu_6Sn_5,2},\ldots$ 

 $\eta_{(Sn),1}, \eta_{(Sn),2}, \dots \eta_{(Sn)i}, \dots$ 

- with  $(\eta_{(Cu),1}, \eta_{(Cu),2}, ..., \eta_{\rho,i}, ...) =$ (1,0,...,0,...), (0,1,...,0,...),...(0,0,...,1,...),...

**Composition field:**  $x_{sn}(\vec{r},t)$ 0

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 $\eta_{(CQ,i)}$ 





### **CALPHAD Gibbs energies**

• COST 531-v3-0 database (+ parabolic extensions)



A. Dinsdale, et al. COST 531-Lead Free Solders: Atlas of Phase Diagrams for Lead-Free Soldering, vols. 1,2 (2008) ESC-Cost office





### IMC-layer growth (1D)







### Comparison with experimental data

*Cu3Sn, T = 180 °C* 



### Parabolic growth constant experiments

Т	k_Cu3Sn	k_Cu6Sn5
150 ℃	0.0010 10 <sup>-6</sup>	0.00032 10 <sup>-6</sup>
180 ℃	0.0032 10 <sup>-6</sup>	0.0038 10 <sup>-6</sup>
200 ℃	0.0043 10 <sup>-6</sup>	0.0071 10 <sup>-6</sup>

*Cu6Sn5, T* = 180 ℃



## Parabolic growth constant in simulations with

$$D_{Sn}^{(Cu)} = 10^{-25} \text{ m}^{2}/\text{s}$$

$$D_{Sn}^{Cu3Sn} = 10^{-15} \text{ m}^{2}/\text{s}$$

$$D_{Sn}^{Cu6Sn5} = 10^{-15} \text{ m}^{2}/\text{s}$$

$$k_{Cu3Sn} = 0.0023 \cdot 10^{-6}$$

$$k_{Cu6Sn5} = 0.0073 \cdot 10^{-6}$$

$$D_{Sn}^{(Sn)} = 10^{-12} \text{ m}^{2}/\text{s}$$

$$J. Janckzak, EMPA$$



Nele Moelans



### Effect of grain boundary diffusion



 $D_{Sn}^{(Cu)} = 2 \cdot 10^{-25}, *2 \cdot 10^{-25} \text{ m}^2/\text{s}$   $D_{Sn}^{Cu3Sn} = 2 \cdot 10^{-15}, *2 \cdot 10^{-13} \text{ m}^2/\text{s}$   $D_{Sn}^{Cu6Sn5} = 2 \cdot 10^{-15}, *2 \cdot 10^{-13} \text{ m}^2/\text{s}$  $D_{Sn}^{(Sn)} = 2 \cdot 10^{-12}, *2 \cdot 10^{-12} \text{ m}^2/\text{s}$ 

 $\gamma_{gb} = 0.25$  J/m<sup>2</sup>









### Growth behavior Cu<sub>3</sub>Sn ?







### **Conclusions**

- What do we need next to improve the quantitative accuracy of phase-field models ?
  - Accurate representation of triple-junction angles outside [100°-140°]
  - CALPHAD Gibbs energies over full composition domain, also for stoichiometric phases and metastable regions
  - Composition and orientation dependent expressions for diffusion and interfacial properties





### Thank you for your attention !

- Acknowledgements
  - Postdoctoral fellow of the Research Foundation Flanders (FWO-Vlaanderen)
  - Simulations were performed on the Flemisch Super Computer (VSC)
  - Projects
    - OT/07/040 (Quantitative phase field modelling of coarsening in leadfree solder joints)
    - IUAP Program DISCO (Dynamical Systems, Control, and Optimization
    - IWT grant SB-73163 (Phase-Field Modelling of the Solidification of Oxidic Systems)

### • More information on http://nele.studentenweb.org





#### WORKSHOP

'Multiscale simulation of heterogeneous materials and coupling of thermodynamic models' January 12-14, 2011

	Home	Program	Registration	Venue	Sponsors	Contact				
50 200	PROGRAM	PROGRAM								
	Confirmed spe	Confirmed speakers:								
	Mathematical	Mathematical multiscale techniques								
	Serge Prudhor	nme (Institute Austin, Goal-or	(Institute for Computational Engineering and Sciences, University of Texas at Austin, USA) Goal-priented adaptivity for multiscale coupling methods							
	Mitchell Luskin (School of Mathematics TBA									
	Frederic Legol	I (CERMI TBA	(CERMICS, ENPC, Paris, France) TBA							
	Marc Geers	(Mecha Netherla TBA	(Mechanical Engineering, Eindhoven University of Technology, the Netherlands) TBA							
	Tim Schulze	(Dept. o Kinetic I	(Dept. of Mathematics, University of Tennesssee, USA) Kinetic Monte Carlo simulation of heterostructured nanocrystalline growth							
	Aleksander Do	nev (Courar Coarse	(Courant Institute, New York University, USA) Coarse-grained particle, continuum and hybrid models for complex fluids							
Webpage:										
	Thermodynamic techniques/Alloys									
http://www.cs.kuleuven.be /conference/multiscale11/	Chris Wolverto	on (Dept. o TBA	of Materials Science a	nd Engineering	, Northwestern Un	iversity, USA)				
	Xavier Gonze	(Physic Louvain TBA	(Physical Chemistry and Physics of Materials, Université Catholique de Louvain, Belgium) TBA							
	Carolyne Cam	p <b>bell</b> (NIST/M Develop Process	(NIST/Metallurgy Division, USA) Development of Multicomponent Diffusion Mobility Databases for Industrial Processing							









### Multi-grain and multi-phase models

- Multi-phase-field model 0
  - Phase fields  $\varphi_1, \varphi_2, \varphi_3, ..., \varphi_p, \quad \sum \varphi_i = 1$
  - **Free energy** •

$$f_{\text{int}} = \sum_{i \neq j} \frac{4\sigma_{i,j}}{\eta_{i,j}} \left\{ \frac{\eta_{i,j}^2}{\pi^2} | \nabla \phi_i \cdot \nabla \phi_j | + \phi_i \phi_j \right\}$$

$$0 < \phi_{i,j} <$$

- Double obstacle, higher order terms, gradient term non-variational
- Interpolation: zero-slope or thermodynamic consistency

Steinbach et al.

MICRESS phase-field code

H. Garcke, B.Nestler, B. Stoth, SIAM J. Appl. Math. 60 (1999) p 295.

L.-Q. Chen and W. Yang, PRB, 50 (1994) p15752

A. Kazaryan et al., PRB, 61 (2000) p14275

- Multi-order parameter models lacksquare

  - Order parameters  $\eta_1, \eta_2, ..., \eta_i(\vec{r}, t), ..., \eta_p, \quad \left(\sum_{i=1}^p \eta_i \neq 1\right)$  61 (200) Interfacial energy  $f_{int} = m \left(\sum_{i=1}^p \left(\frac{\eta_i^4}{4} \frac{\eta_i^2}{2}\right) + \sum_{i=1}^p \sum_{j < i}^p \gamma_{i,j} \eta_i^2 \eta_j^2 + \frac{1}{4}\right) + \frac{\kappa(\eta)}{2} \sum_{i=1}^p \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{i=1}^p \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{i=1}^p \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta_j^2 \eta_j^2 \eta_j^2 + \frac{1}{4} + \frac{\kappa(\eta)}{2} \sum_{j < i}^p \eta_j \eta_j^2 \eta$





## Multi-grain and multi-phase models

- Vector valued model
  - Orientation field ( $\phi$ ) and phase field ( $\phi$ )
  - Free energy  $f_{\text{int}} = f(\phi, |\nabla \phi|, |\nabla \theta|)$
- 2-phase solidification
  - Phase fields

$$\varphi_1, \varphi_2, \varphi_3, \quad \sum_{i=1}^{3} \varphi_i = 1$$

- Fifth order interpolation functions  $g_i(\phi_1, \phi_2, \phi_3)$ 
  - Zero-slope and thermodynamic consistent
  - Order  $g_i$  increases with number of phase-fields
- Multi-order parameter + 4th order gradient terms
- I.M. McKenna, M.P. Gururajan, P.W. Voorhees, J. Mater. Sci., 44 (2009) p2206
- Phase field crystal and amplitude equations

Nele Moelans Third annual workshop HERO-M, Saltsjöbaden, Sweden, May 17-18, 2010



*R. Kobayashi, J.A. Warren, W.C. Carter, Physica D, 119* (1998) p415

R. Folch and M. Plapp, PRE, 72 (2005) n ° 011602



### **Bounding Box Algorithm**







- Initialization by random nucleation
  - Generate sphere-shaped grain uniformly over microstructure
  - Generate particles uniformly over microstructure
- Set-up sparse data structure
  - Determine bounding box for every grain
  - Create object for every grain

